

Teaching Physics with Physmatics

Enseñando Física con Fismática

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Quantum TimeLord Virtual Academy

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**XXXVIII Biennial of Physics of the Spanish Royal Physics
Society (R.S.E.F.) (Murcia, 11-15 July 2022)**

**XXXVIII Reunión Bienal de la Real Sociedad Española de Física
(Murcia, del 11 al 15 de julio de 2022), Earth planet
Milky Way Galaxy, Laniakea, Known Universe (The Multiverse)**

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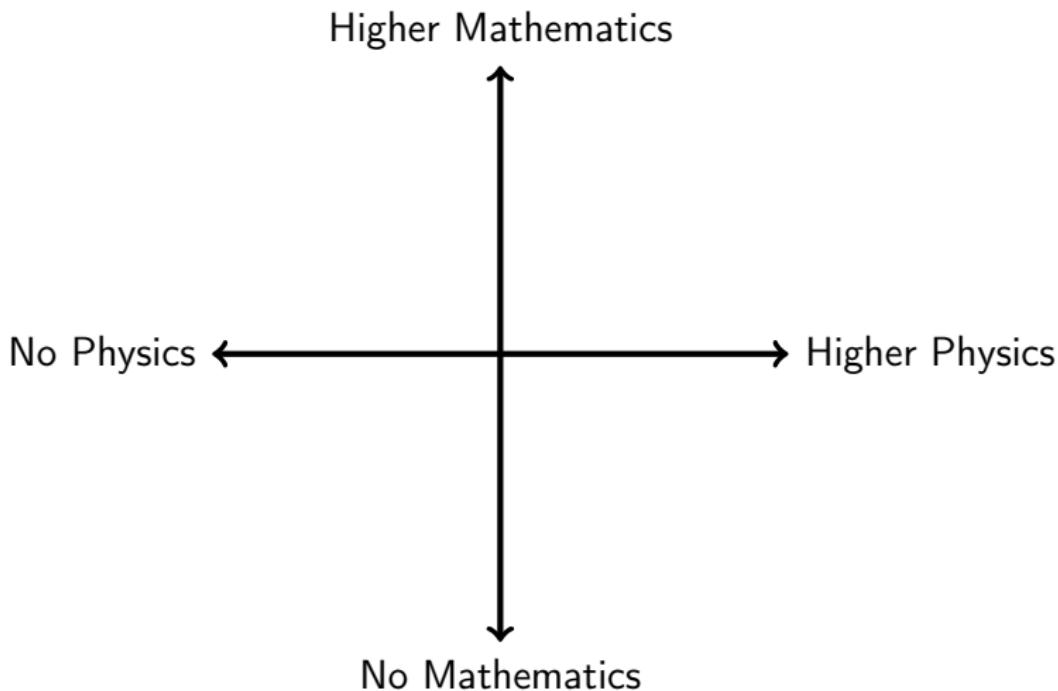
What Physmatics IS/Lo que es la Fismática

Gráfica de lo que es la Ciencia de la Fismática en comparación a otros enfoques/Plot of what Physmatics is compared to other approaches. . .

No Physics ← → Higher Physics

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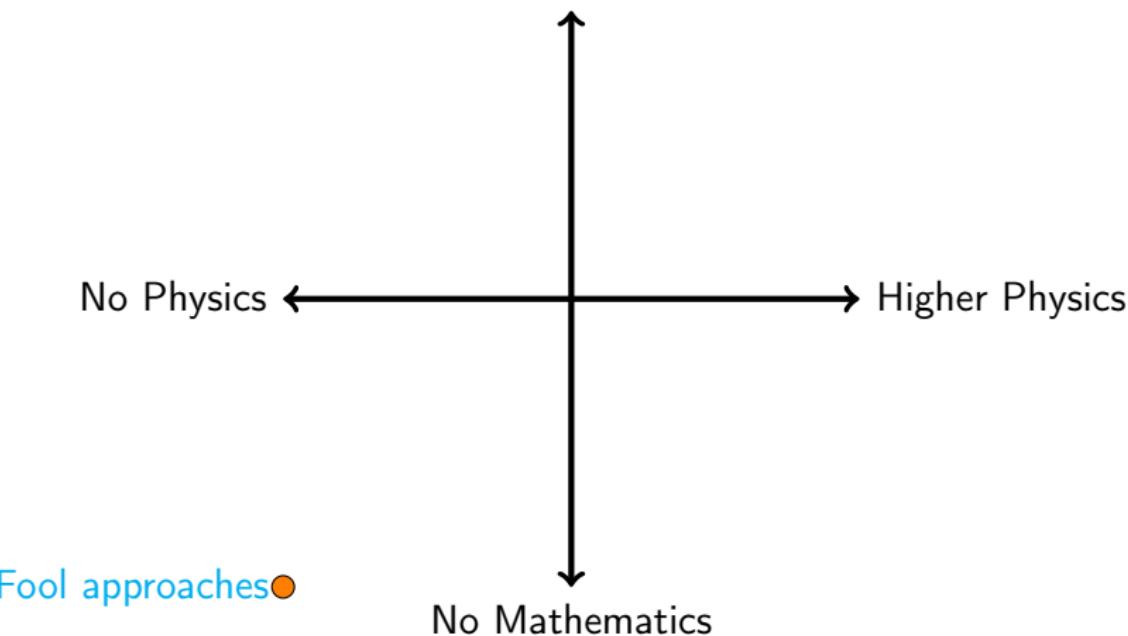
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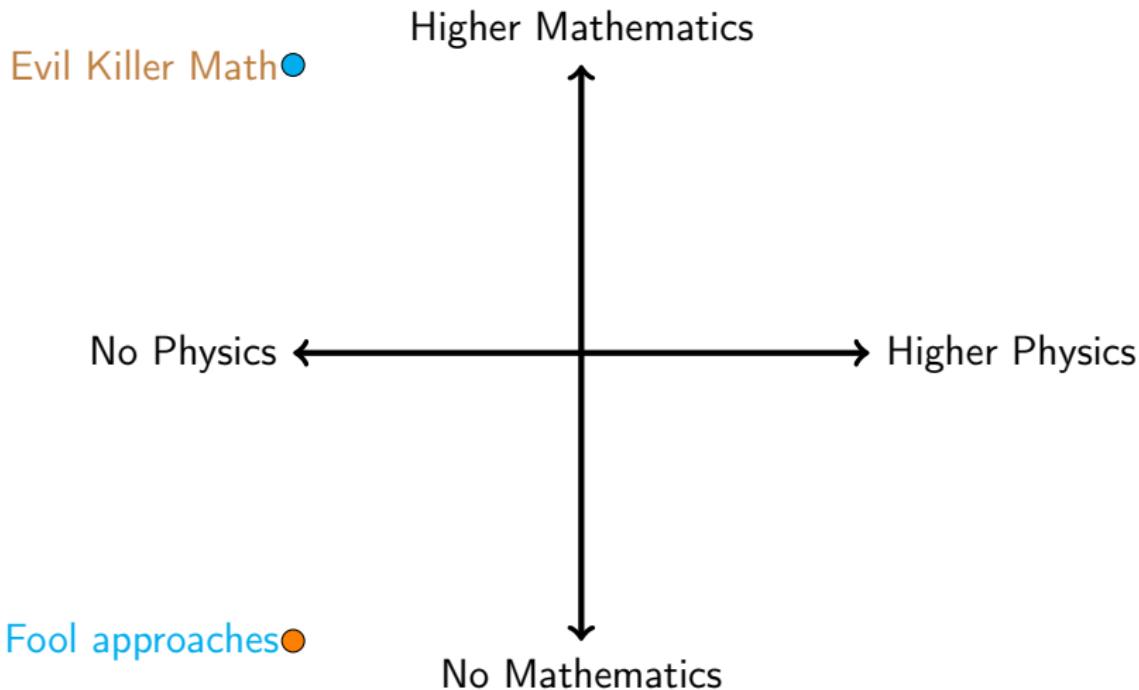
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Higher Mathematics



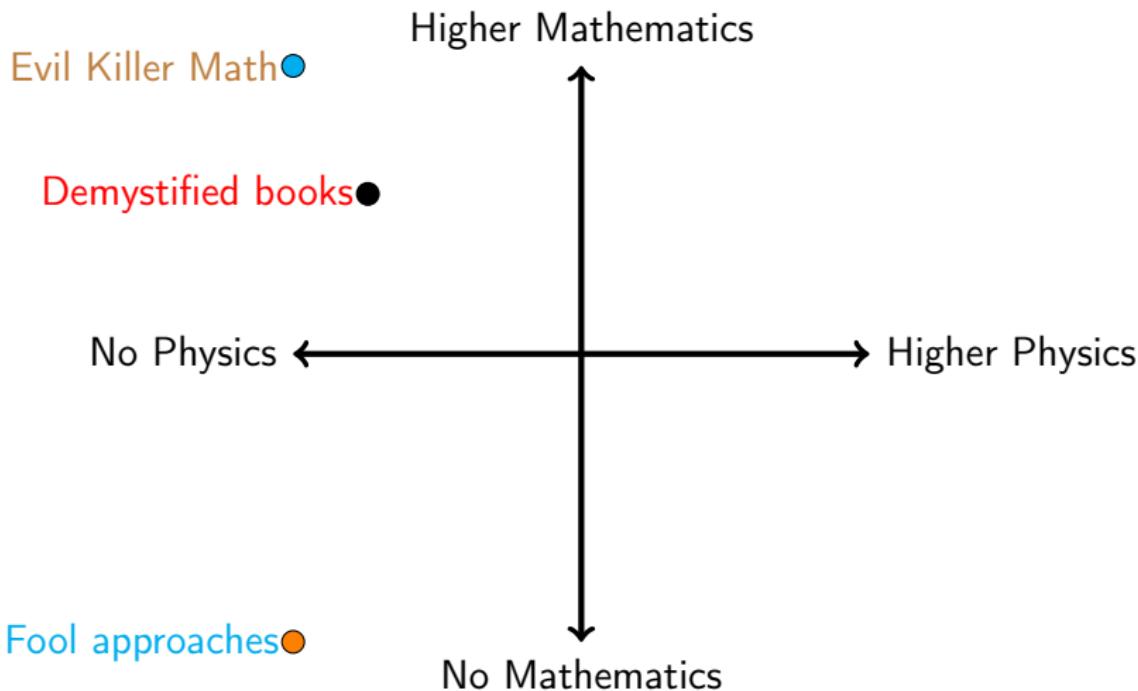
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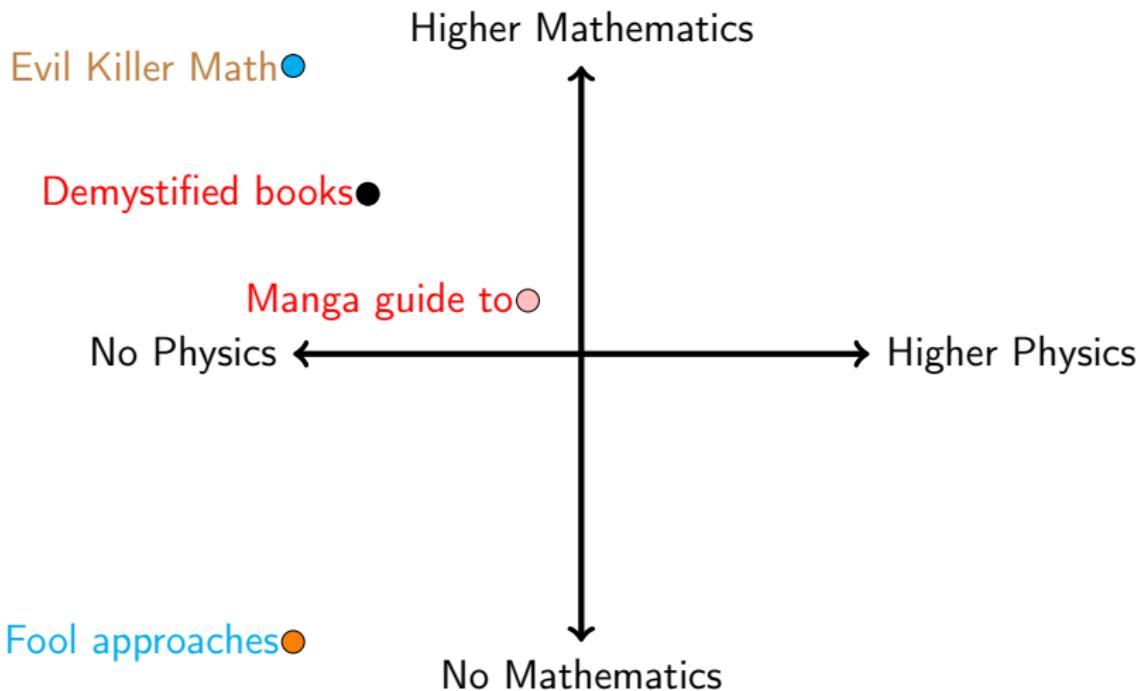
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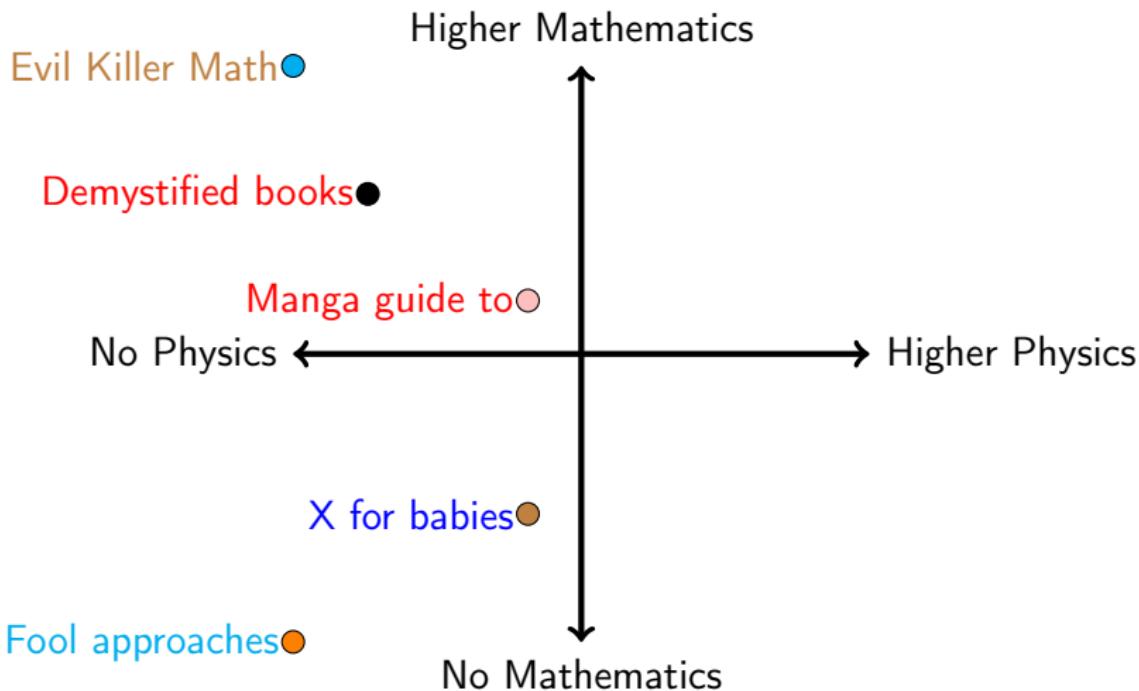
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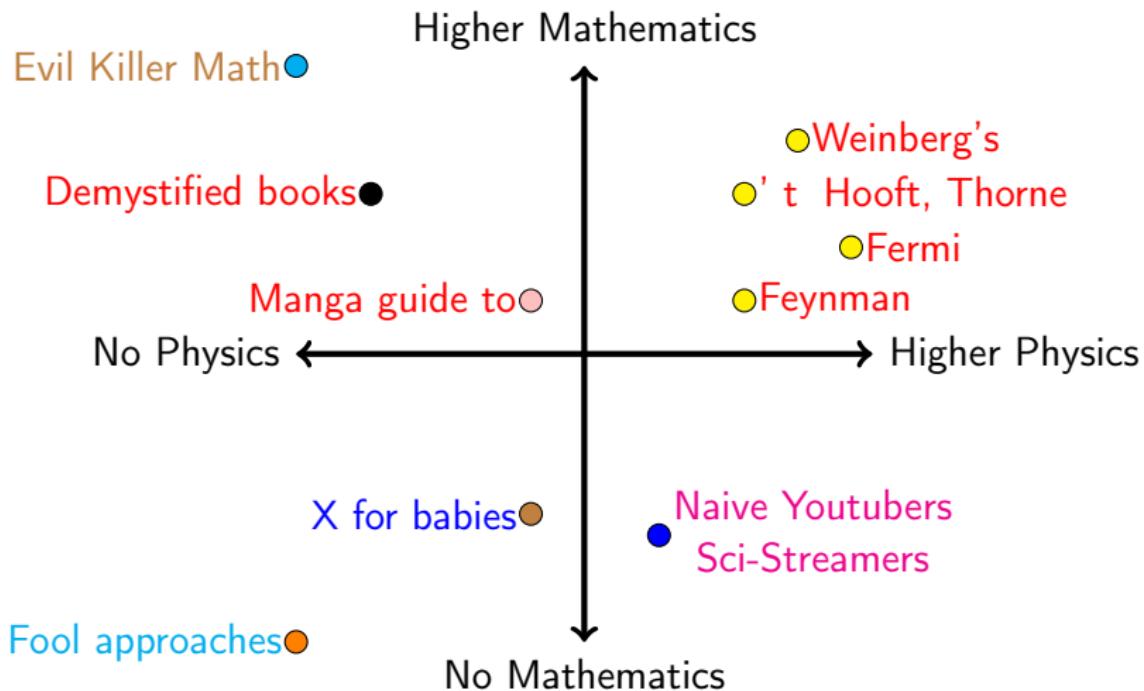
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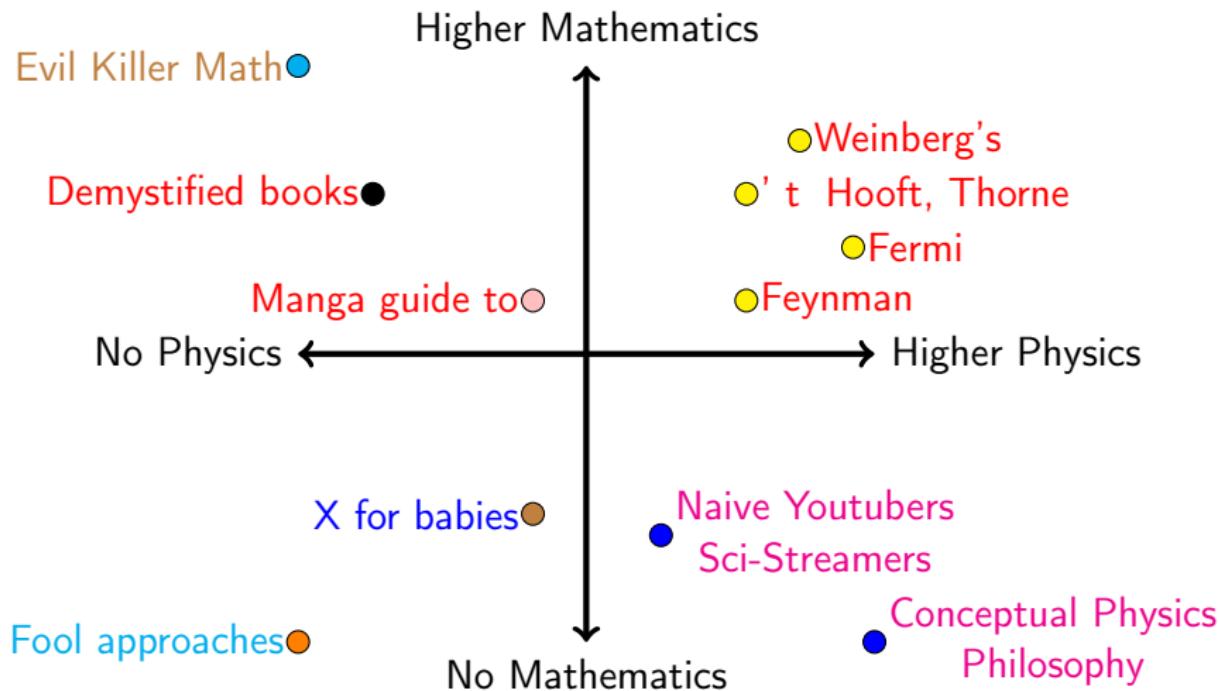
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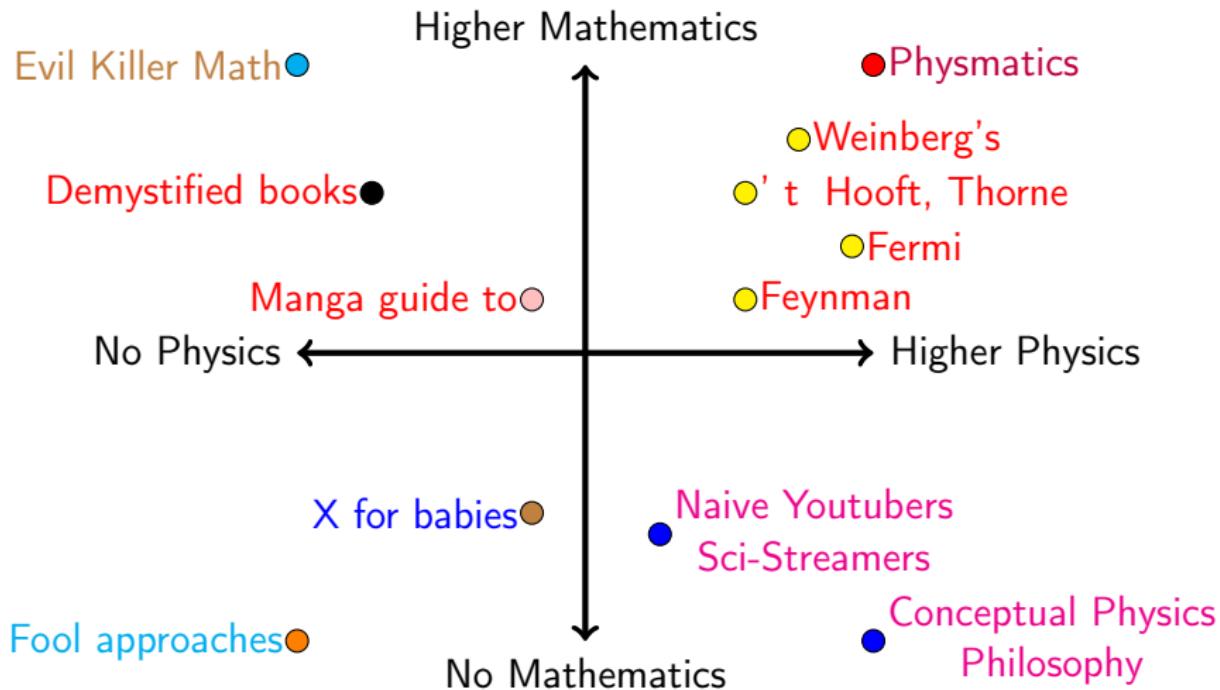
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The hypersphere(\mathbb{I})/la hiperesfera(\mathbb{I})

Hypersphere

Hypersphere (euclidean) or n -sphere is the geometrical locus or manifold S^n with $(n-1)$ -dimensional constraint $\sum_{i=1}^n x_i^2 = x_1^2 + \cdots + x_n^2 = R^2$.

Hypervolume and hypersurface for n -spheres

The hypervolume $V(S^n)$ and hypersurface $\Sigma(S^n)$ is calculated as follows:

$$V_n = \frac{\Gamma(1/2)^n R^n}{\Gamma\left(\frac{n}{2} + 1\right)} \quad \Sigma_{n-1} = \frac{dV_n}{dR} = \frac{n\Gamma^n(1/2)R^{n-1}}{\Gamma\left(\frac{n}{2} + 1\right)} \quad (1)$$

where $\Gamma(1/2) = \sqrt{\pi} = (-1/2)!$. Remark: $V(S^\infty) = 0$, and the volume of the 23-sphere unit sphere is equal to the Leech lattice volume

$\Lambda_{24} = \pi^{12}/12!$ behind the symmetry of the monster group M . Dimensional recurrence: $V_n = \frac{R\Sigma_{n-1}}{n}$.

The hypersphere/La hiperesfera(I)

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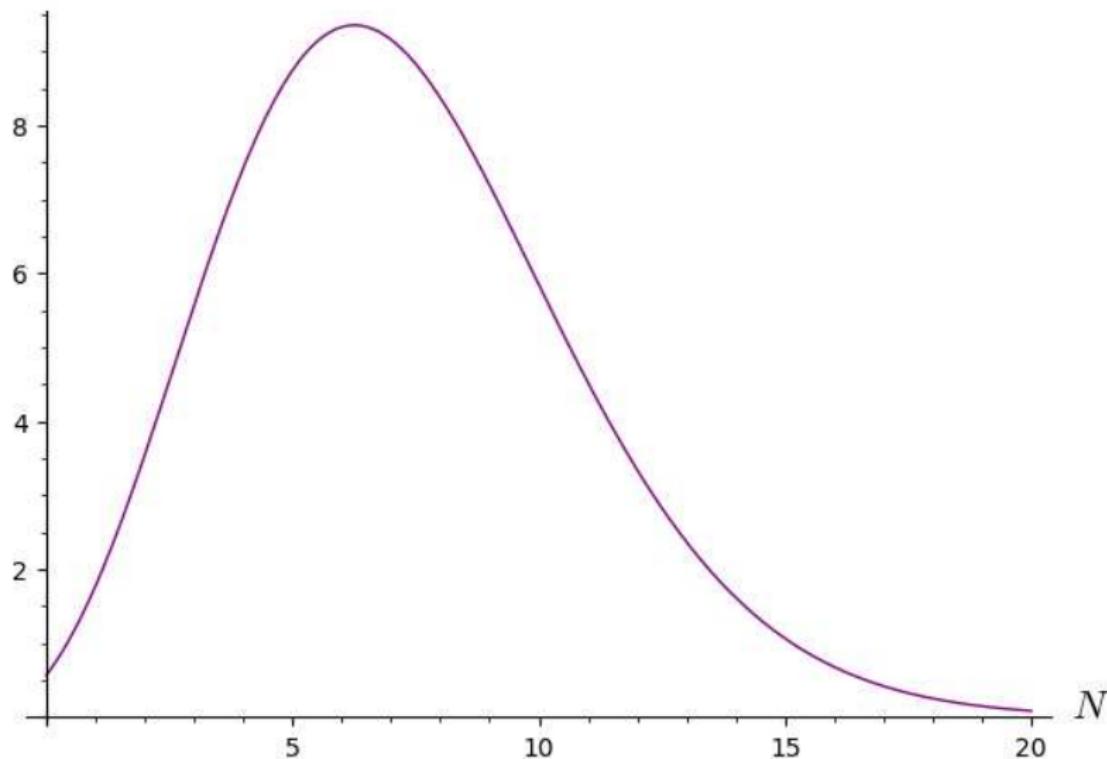
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$\Lambda_{24} = \pi^{12}/12!$ behind the symmetry of the monster group M . Dimensional recurrence: $V_n = \frac{\Sigma_{n-1} R}{n}$, $V(N) = 2\pi R^2 V(N-2)/N$.

Gráfica/Graphics

$V(N)$ Unit hypersphere volume, volumen de hiperesfera unidad



Special functions/Funciones especiales

Solve, exactly, the following equations/**Resolver, exactamente, las siguientes ecuaciones:**

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- $xa^x = Y$
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- $\ln(A + BX) + CX = \ln(D)$
- $x^{x^{x^{\dots}}} = y$
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Hint: Use Lambert W-function! Should we teach it at (high)-school/University?

Pista: Usar la función W de Lambert. ¿Deberíamos enseñarla en la Universidad o en el IES?

Deformed calculus/Cálculo deformado

Everyone knows what a derivative is...

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and in the space of analytic functions if $\gamma/\alpha \in \mathbb{Z}$

$$D_{p,q}^{\alpha,\beta,\gamma}f(z) = \frac{f(p^{-\alpha}z)p^{-\beta} - f(q^\alpha)q^\beta}{(p^{-\gamma} - q^\gamma)z^{\gamma/\alpha}} \quad (4)$$

Moreover, take the fractional Riemann-Liouville derivative

$${}_aD_x^{-\alpha}f(x) = \frac{1}{\Gamma(\alpha)} \int_a^x f(t)(x-t)^{\alpha-1} dt. \quad (5)$$

There are more, but...

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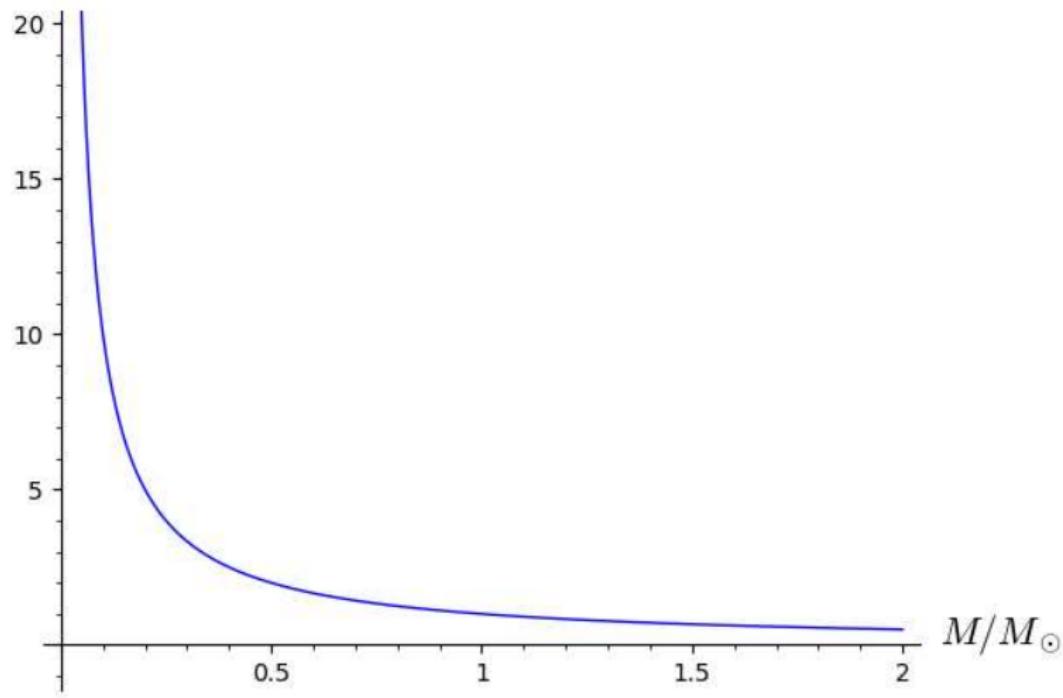
Proportion and dimensional analysis

- A tool for Black hole chemistry? ¿Un instrumento para la Química de agujeros negros?
- Via power laws/Via leyes de potencias . .

Proportion and dimensional analysis

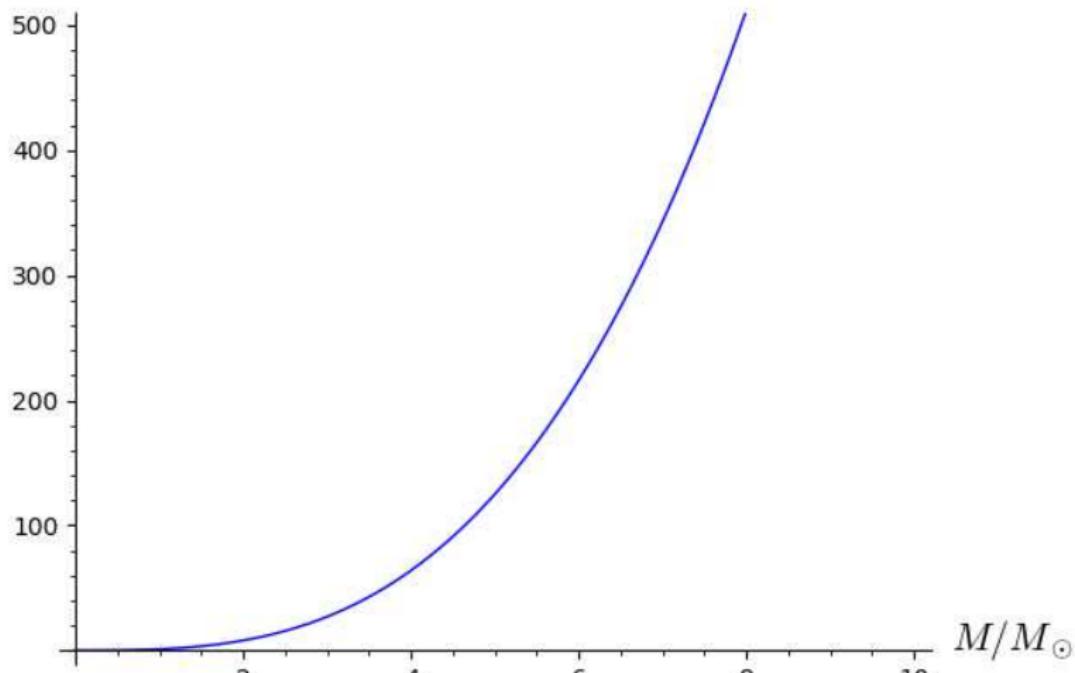
$$\frac{8\pi G k_B T}{\hbar c^3}$$

Black hole temperature, in certain units $T = \frac{\kappa_q}{M}$

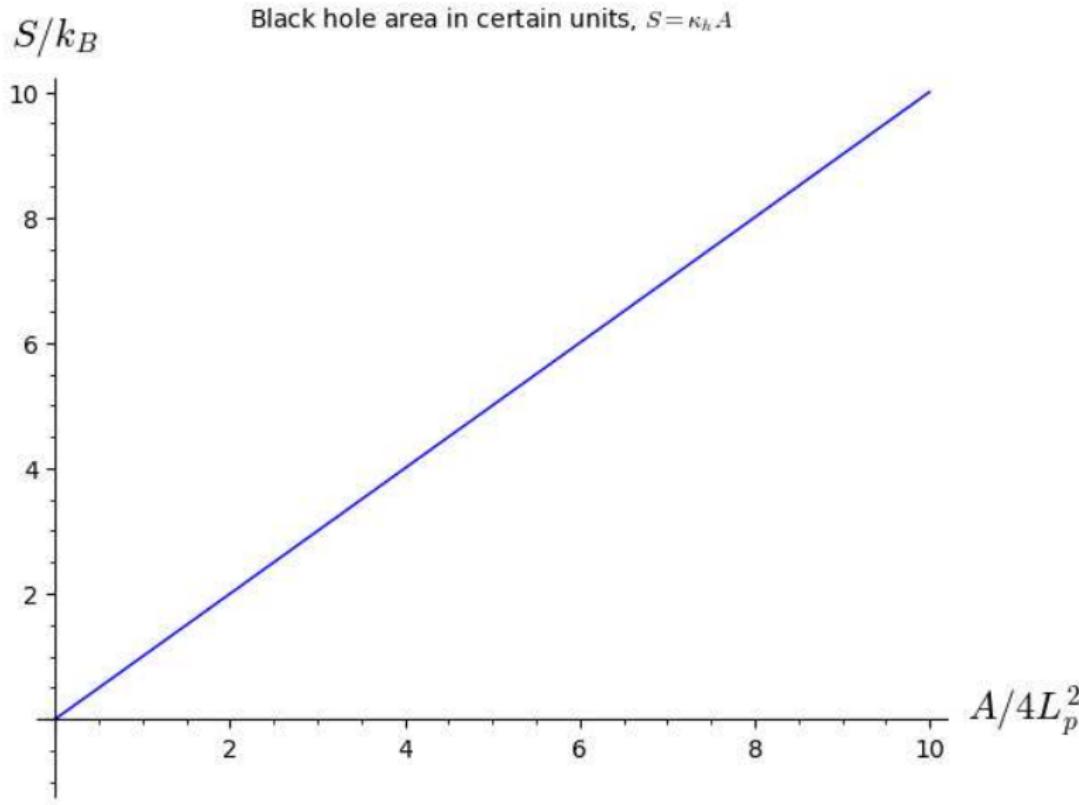


Proportion and dimensional analysis

$$\frac{\hbar c^4 \alpha_e t_e}{5120\pi G^2}$$
 Black hole evaporation time in certain units, $t_e = K M^3$



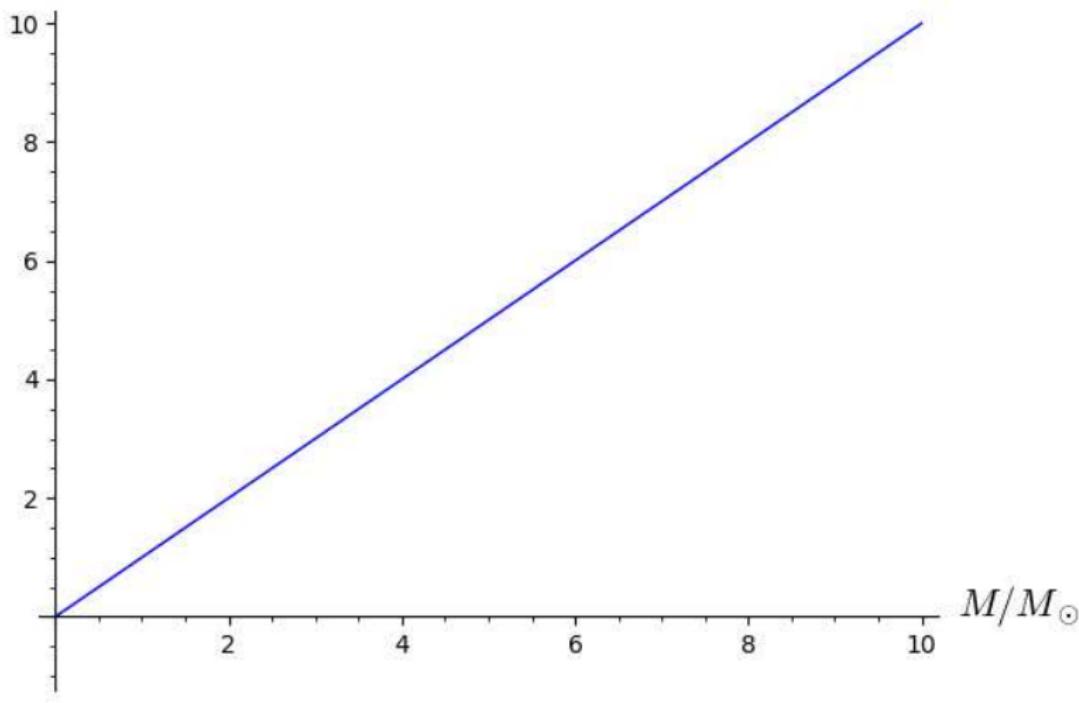
Proportion and dimensional analysis



Proportion and dimensional analysis

$$\frac{c^3 t_S}{\pi G}$$

Time to singularity in certain units, $t_S = \kappa M$



Kepler 3rd law variants(I)

Usual Kepler 3rd law: $T^2 = \frac{4\pi^2}{GM_B} R^3$. What about some variants? For Kerr black holes:

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Usual Kepler 3rd law: $T^2 = \frac{4\pi^2}{GM_B} R^3$. What about some variants? For Kerr black holes:

$$\Omega = \pm \frac{M^{1/2}}{r^{3/2} \pm aM^{1/2}} \quad (6)$$

where $M = G_N m$ is the mass in gravitational natural units, a is the Kerr rotation parameter $a = cJ/M = J/Mc$. In terms of complete dimensional constants reads

$$\Omega = \pm \frac{\sqrt{G_N M}}{r^{3/2} \pm \chi \left(\frac{\sqrt{G_N M}}{c} \right)^3} = \pm \frac{c^3}{GM} \left(\pm \chi + \left(\frac{c^2 r}{G_N M} \right)^{3/2} \right)^{-1} \quad (7)$$

Kepler 3rd law variants(II): beyond standard gravity

For gravitational theories with effective potential:

$$V_e = -\frac{GM}{r} \left(1 + A \frac{M^p}{r^p}\right) + \frac{L^2}{2\mu^2 r^2}$$

circular orbit condition reads $V' = 0$, and with $L = \mu r^2 \Omega^2$ you get the generalized Kepler 3rd law

$$\Omega^2 = \frac{GM}{r^3} \left(1 + \frac{M^p A(p+1)}{r^p}\right) \quad (8)$$

More? Take the Finslerian-like 3rd law modification:

$$\frac{r^3}{T^2} = \left(1 - \frac{A(r)}{r^4}\right) \frac{GM}{4\pi^2} \quad (9)$$

Kepler 3rd law variants(II): strong gravity modified QG?

Recently, it has been proposed a modified gravitational law with effective potential energy:

$$U_e = -\frac{GMm}{r} - \lambda Mm \ln \left(\frac{r}{r_0} \right)$$

giving

$$F(r) = -\frac{GMm}{r^2} - \frac{\lambda Mm}{r} = -G_e \frac{Mm}{r^2} \quad (10)$$

where you get an effective gravitational constant

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Similar ideas are proposed by asymptotically safe gravity approaches (by Weinberg and others), where $G_e = G(r)$, or even superstrings with $G = G(r, t) = g_s^2 L_s^2 e^{\phi(r, t)}$.

Quantum gravity, string theory and molecular forces

We can practice the usual $F = -\nabla E_p$ between conservative fields/forces and gravity. For instance, General Relativity plus Quantum Gravity at one loop corrections (effective theory) provide the potential energy

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$$E_p = U = -\frac{GM_1 M_2}{r} \left(1 + a \frac{G(M_1 + M_2)}{c^2 r} + b \frac{G\hbar}{c^3 r^2} \right) \quad (12)$$

with $a = 3$ (GR) and $b = 41/10\pi$ (QG at one loop). Reciprocally, we could take the effective force in 26d bosonic string theory, namely

$$F_N = G_{26d} \frac{Mm}{r^{24}} \quad (13)$$

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Exercise: calculate U_{26d} !. It can be shown $G_{26d} = g_s^2 L_s^{24}$ in certain units. Intermolecular force between diatomic molecules can be approximated by the central force:

$$f(r) = -\frac{K_1}{r^6} + \frac{K_2}{r^{12}}$$

Find out the potential energy for this force.

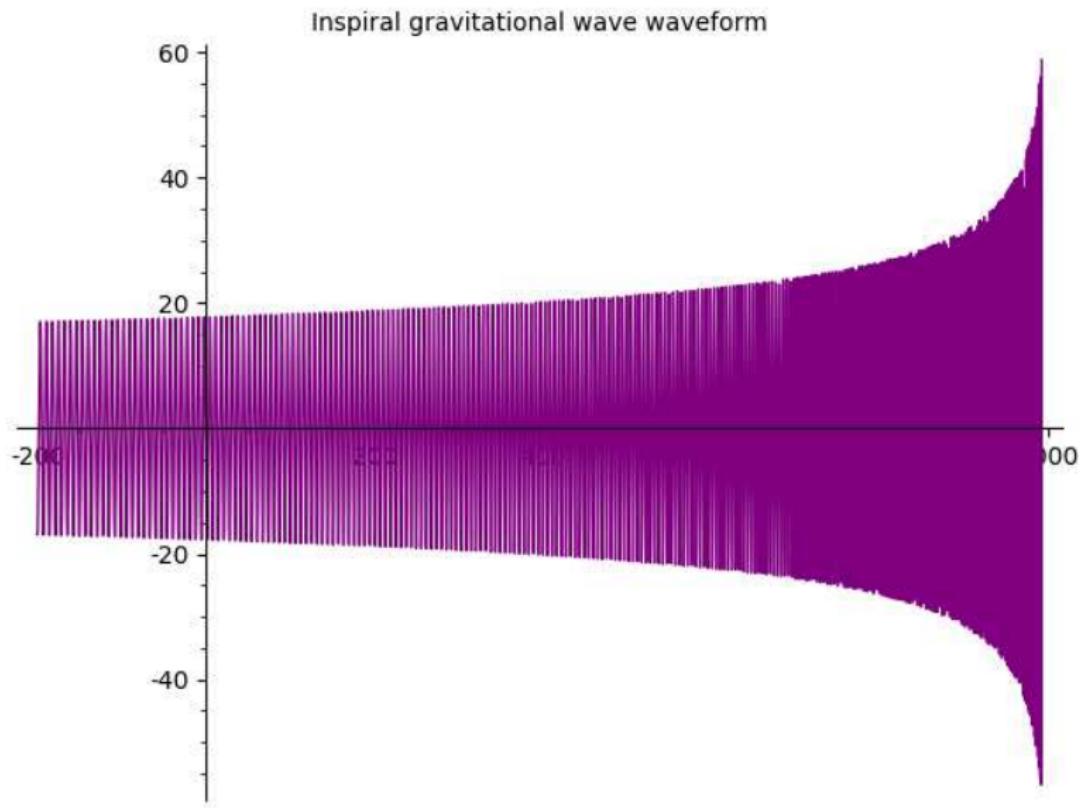
Waves(I): inspiral GW

Certain oscillating system moves with waveform, S.I. units,:

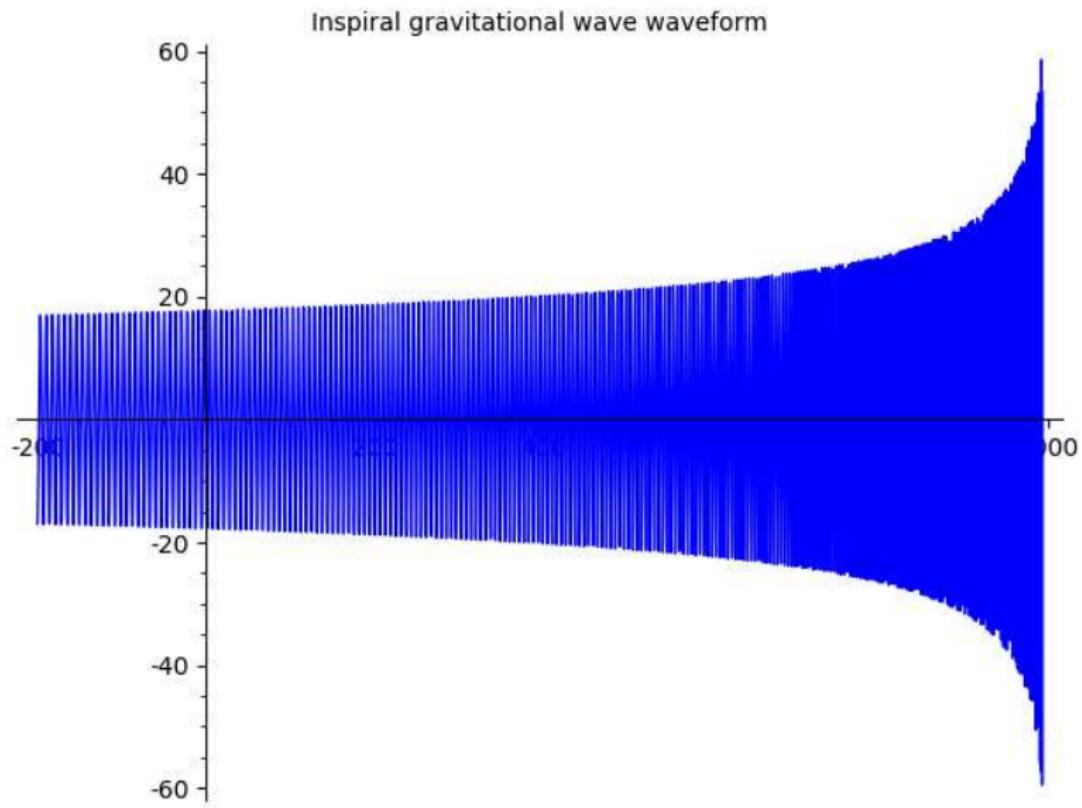
$$\Psi = 100(1000 - t)^{-1/4} \cos \left((10(1000 - t))^{5/8} \right)$$

- a) Is this a SHO? Explain the answer.
- b) Plot $y(t)$.
- c) Calculate the vibrational speed and acceleration, $v = d\Psi/dt$, $a = d^2\Psi/dt^2$.
- d) Calculate when is maximum the vibrational speed and its value, and the value of Ψ at that value. Comment the results.

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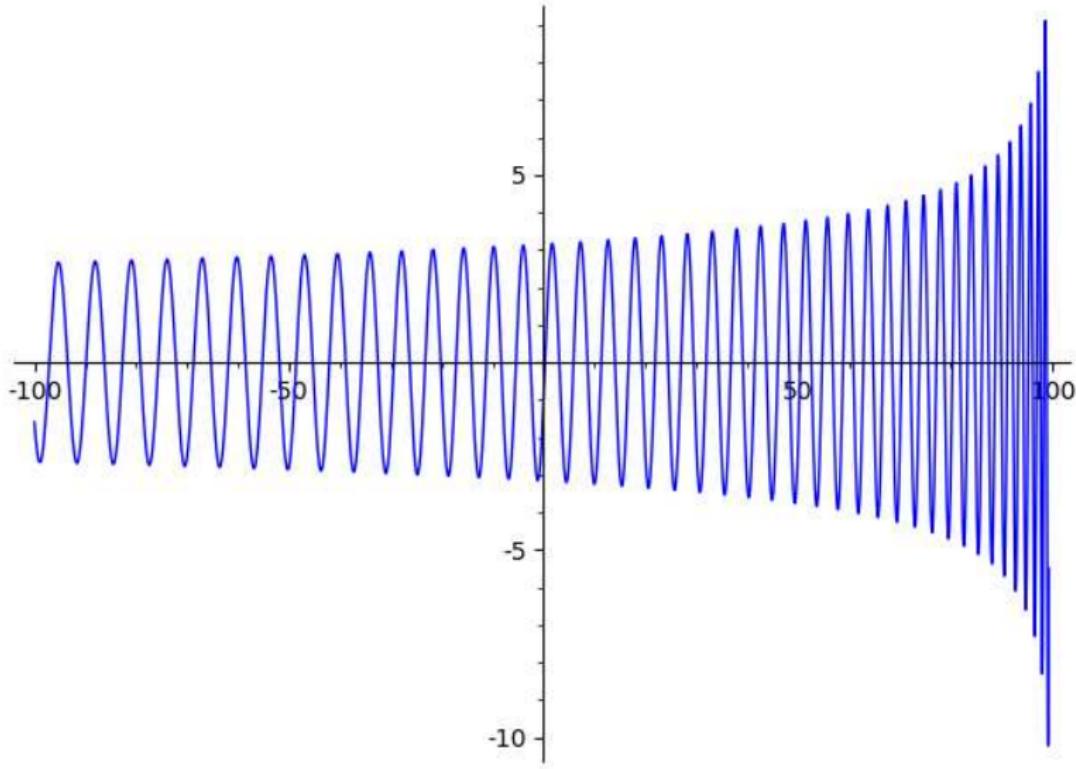


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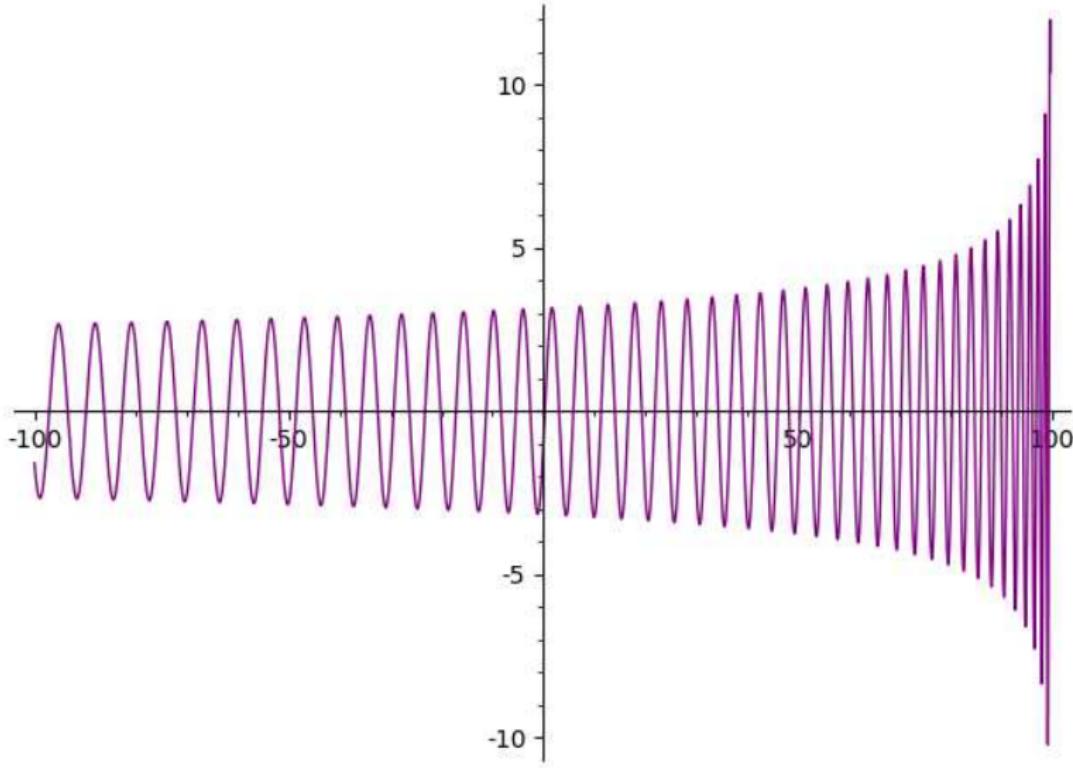
Waves(I): inspiral GW

Inspiral gravitational wave waveform scaled and zoomed



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Waves(II): Axion DM model

A model for dark matter [5], which density reads $\rho_{DM} = 0,4 \text{ GeV/cm}^3$, predicts a distribution with waveform due to ultralight (pseudo)scalar particles. Scalar waves along positive OX are given by:

$$\phi(x, t) = 2 \frac{\sqrt{2\rho_{DM}}}{m_\phi c^2} \cos\left(2\pi \frac{m_\phi c x}{h}\right) \sin\left(2\pi \frac{m_\phi c^2 t}{h}\right)$$

with $c = 3 \cdot 10^8 \text{ m/s}$ is the speed of light, $h = 6,63 \cdot 10^{-34} \text{ J} \cdot \text{s}$ is the Planck constant, m_ϕ and is the mass of particles ϕ . Determine, if $E_\phi = m_\phi c^2 = 10 \cdot 10^{-22} \text{ eV} = 1 \text{ zeV}$:

- Type of wave and explanation, the propagation speed v_p , wavelength λ_ϕ , frequency f_ϕ and period T_ϕ . (0.5 points=0.1x5).
- Amplitudes for ϕ and ϕ^2 with S.I. units (0.5 points).

Waves(II): Axion DM model

- c) The number of maxima we would expect to see in the square of particle distribution $\phi(0, t)^2$, using a haloscopic detector, in a time equal to the period, $t = T_\phi$. (0.5 points)
- d) Number of particles per cubic meter if all the dark matter is locally made of ϕ particles. What happens to the field and the wavefunction in the limit of massless field, i.e. $m_\phi = 0$? (0.5 points)

Data: $e = 1,6 \cdot 10^{-19} C$

Cosmography(I)

In Cosmography we define, for $a(t)$:

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$$\text{Hubble function : } H(t) = \frac{1}{a} \frac{da}{dt} \quad (14)$$

$$\text{deceleration function : } q(t) = -\frac{1}{aH^2} \frac{d^2a}{dt^2} \quad (15)$$

$$\text{jerk : } j(t) = \frac{1}{aH^3} \frac{d^3a}{dt^3} \quad (16)$$

$$\text{snap : } s(t) = \frac{1}{aH^4} \frac{d^4a}{dt^4} \quad (16)$$

$$\text{crackle/lerk function : } l(t) = \frac{1}{aH^5} \frac{d^5a}{dt^5} \quad (17)$$

Cosmography(II)

Cosmographic parameters in current time are denoted by $(H_0, q_0, j_0, s_0, l_0)$ and are a target of present and future cosmological measurements. Since:

$$a(t) = a(0) + \frac{da}{dt}(0)t + \frac{d^2a}{dt^2}(0)t^2 + \cdots + \frac{1}{n!} \frac{d^n a}{dt^n}(0)t^n + \mathcal{O}(t^{n+1}) \quad (18)$$

and then H with redshift z

$$H(z) = H(0) + \frac{dH}{dz}(0)z + \frac{d^2H}{dz^2}(0)z^2 + \cdots + \frac{1}{n!} \frac{d^n H}{dt^n}(0)z^n + \mathcal{O}(z^{n+1}) \quad (19)$$

a) Check, with explicit calculations, the relationships below :

$$\dot{H} = -H^2(1 + q)$$

$$\ddot{H} = H^3(3q + j + 2), \quad \ddot{\dot{H}} = H^4(-3q^2 - 12q - 4j + s - 6)$$

$$\ddot{\dot{\dot{H}}} = H^5(30q^2 + 60q + 10qj + 20j - 5s + l + 24)$$

b) Find out the dimensions and units of cosmographic parameters, explaining why the Hubble law is related to them, $v = H(z)d(z)$.

Cyclic model of the universe

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In the talk and article [1], cosmic equation for a Universe dominated by a network of “domain walls” is proposed and written, and solved as a cyclic model:

$$\ddot{a} + \frac{|\Lambda|}{3}a = \frac{4\pi G_N c}{3}, \text{ with solution } a(t) = \frac{1}{\gamma\omega} \left(1 + \sqrt{1-\gamma} \cos(\omega t)\right) \quad (20)$$

where the frequency and width of oscillation are defined as:

$$\omega = \sqrt{\frac{|\Lambda|}{3}}, \quad \gamma = \frac{3|\Lambda|}{(4\pi G_N c)^2} \sim \frac{a_-}{a_+} \quad (21)$$

We used the Friedmann equations for FRW and the state equation $p = \omega\rho$, with $\rho = -2/3$ (domain walls).

Inflation and the Multiverse (I)

How many Universes in the Multiverse? [3], i.e., ¿cuántos universos hay en el Multiverso? Respuestas posibles (A.Linde et al.; cf. ideas de A. Vilenkin y A. Guth):

- Slow-roll inflation provides $\mathcal{N}_{\text{efolds}} \sim e^{e^{3N}}$. If $N = 60$ (usual hypothesis), then $\mathcal{N} \sim e^{e^{180}} \sim 10^{10^{77}} \sim 10^{S_{BH}(1M_\odot)}$.
- (Chaotic) Eternal inflation provides, $N = cS_{dS} \sim c/m$:
$$\mathcal{N} \sim e^{e^{3N}} \sim e^{e^{3c/m}} \sim 10^{10^{10^7}} \gg 10^{\text{googol}}$$
- Cosmological constant universes, Λ non-zero, imply
$$S_p \approx H \frac{1+3\omega}{1+\omega} |\Lambda|^{-\frac{1+3\omega}{2+2\omega}} = H^{3/2} |\Lambda|^{-3/4} \rightarrow \mathcal{N} \sim 10^{10^{82}}.$$
- Universes in the string landscape (M is the number of dS vacua):

$$\mathcal{N} \approx \sum_{j=1}^M \exp(|\Lambda|^{-3/4}) \sim e^{0.75M}, \quad \text{Popularly } M \sim 10^{500}, \mathcal{N} \sim 10^{10^{375}}$$

Other options: $M \sim 10^{272000}$ (F-theory), $M \sim 10^{15}$, $\Lambda \sim 10^{-122}$, $N_{\max} \sim \log|\Lambda| \sim 290$, and $N(\text{efolds}) \sim 70$.

Inflation and the Multiverse (II)

- Our (observable) Universe has an entropy bound $S_{dS} \leq |\Lambda|^{-3/4} \sim 10^{90}$.
- Milky way black hole entropy: $S_{MWBH} \sim 10^{100} \sim \text{googol}$.
- The number of observers with masses about $M \sim 10^2 \text{ kg}$, and height 1m is bound by Bekenstein formula

$$\mathcal{N}_{obs} \leq S_{Bek} = e^{2\pi MR} \leq e^{10^{45}}$$

Remarkly, the number of an intelligent brain observer is about $\mathcal{N} \sim 10^{10^{16}} >> \mathcal{N}_{dS-Vacua}$, a brain seems to have more configurations than the expected possible geometries of the Universe(Multiverse).

Inflation and the Multiverse (II)

- Our (observable) Universe has an entropy bound $S_{dS} \leq |\Lambda|^{-3/4} \sim 10^{90}$.
- Milky way black hole entropy: $S_{MWBH} \sim 10^{100} \sim \text{googol}$.
- The number of observers with masses about $M \sim 10^2 \text{ kg}$, and height 1m is bound by Bekenstein formula

$$\mathcal{N}_{obs} \leq S_{Bek} = e^{2\pi MR} \leq e^{10^{45}}$$

Remarkly, the number of an intelligent brain observer is about $\mathcal{N} \sim 10^{10^{16}} >> \mathcal{N}_{ds-Vacua}$, a brain seems to have more configurations than the expected possible geometries of the Universe(Multiverse).

En la inflación, un campo escalar en modelos de inflación cáratica eterna $V(\phi) = 0,5m^2\phi^2$ tiene un comportamiento de MAS modulado en amplitud

$$\phi(t) = \Phi(t) \cdot \sin(m\phi), \text{ with } \Phi(t) = \frac{M_p}{\sqrt{3\pi}mt} \sim \frac{M_p}{2\pi\sqrt{3\pi}N} \quad (22)$$

and N is the number of oscillations since the end of inflation!



Inflation energy scale and tensor-to-scalar ratio

General inflation models use ODE like the Lame or Mathieu equations (H is a friction term), or elliptic functions. Moreover, the energy scale of inflation is related to the tensor-to-scalar perturbation ratio.

Inflation energy scale

$$V^{1/4} \approx \left(\frac{3\pi^2}{2} r \mathcal{P}_s \right)^{1/4} M_P = \left(\frac{r}{0,01} \right)^{1/4} \cdot 1,06 \cdot 10^{16} \text{GeV} \quad (23)$$

Primordial gravitational waves triggered by inflation in the Early Universe (even in the beginning of time or “before” the Big Bang) are a hot target of current and future research!

References for the last 2 slides:

- ① *Primordial Gravitational Waves from Cosmic Inflation.* Mike S. Wang. Mathematical Tripos Part III Essay 75 (colour in electronic version)
Submitted 5th May 2017, updated 26th August 2017.
- ② *Towards the Theory of Reheating After Inflation.* Lev Kofman, Andrei Linde, Alexei A. Starobinsky.

Contenido

1 What is Phsmatics? ¿Qué es la Fismática?

2 Teaching with Phsmatics/Enseñando con Fismática

3 Bibliography

References/Referencias, Bibliography/Bibliografía

- [1] *A Simple Harmonic Universe*, P. Graham, B. Horn, S. Rajendran and G. Torroba. ArXiv: <https://arxiv.org/abs/1109.0282v2>
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- [3] *How many universes are in the multiverse?*, Andrei Linde and Vitaly Vanchurin. ArXiV:<http://arxiv.org/abs/0910.1589v3>
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- [5] *Dark matter or strong gravity?*, Saurya Das and Sourav Sur, arXiv: <https://arxiv.org/abs/2205.07153>
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- [7] *EPFL Lectures on General Relativity as a Quantum Field Theory*, John F. Donoghue , Mikhail M. Ivanov, and Andrey Shkerin. ArXiv: <https://arxiv.org/abs/1702.00319>.
- [8] *Mass and Motion in General Relativity*, Springer. Luc Blanchet, Alessandro Spallicci, Bernard Whiting.

I could do this every day of my life...



My passion is Physmatics...Do I ever sleep?

THAT'S MY SECRET

I NEVER SLEEP



Please... I am grateful for your attention!



I know... Things just got out of hands... I apologize!



Gracias/Thank you!



Back-up slides

Back-up slides

Objetos geométricos: la hiperesfera(II)

$$V(S^0) = 2R \quad (24)$$

$$V(S^1) = \pi R^2 \approx 3,14159 R^2 \quad (25)$$

$$V(S^2) = \frac{4}{3} \pi R^3 \approx 4,11879 \quad (26)$$

$$V(S^3) = \frac{\pi^2}{2} R^4 \approx 4,9348 R^4 \quad (27)$$

$$V(S^4) = \frac{8\pi^2}{15} R^5 \approx 5,26379 R^5 \quad (28)$$

The amazing vanishing sphere volume with increasing dimensions!!!!!!!

Objetos geométricos: la hiperesfera(II)

$$V(S^5) = \frac{\pi^3}{6} R^6 \approx 5,16771 R^6 \quad (24)$$

$$V(S^6) = \frac{16\pi^3}{105} R^7 \approx 4,72477 R^7 \quad (25)$$

$$V(S^7) = \frac{\pi^4}{24} R^8 \approx 4,05871 R^8 \quad (26)$$

$$V(S^8) = \frac{32\pi^4}{945} R^9 \approx 3,29851 R^9 \quad (27)$$

$$V(S^9) = \frac{\pi^5}{120} R^{10} \approx 2,55016 R^{10} \quad (28)$$

The amazing vanishing sphere volume with increasing dimensions!!!!!!



Objetos geométricos: la hiperesfera(II)

$$V(S^{10}) = \frac{64\pi^5}{10395} R^{11} \approx 1,8841 R^{11} \quad (24)$$

$$V(S^{11}) = \frac{\pi^6}{720} R^{12} \approx 1,33526 R^{12} \quad (25)$$

$$V(S^{12}) = \frac{128\pi^6}{135135} R^{13} \approx 0,919629 R^{13} \quad (26)$$

$$V(S^{13}) = \frac{\pi^7}{5040} R^{14} \approx 0,599265 R^{14} \quad (27)$$

$$V(S^{14}) = \frac{256\pi^7}{2027025} R^{15} \approx 0,381443 R^{15} \quad (28)$$

The amazing vanishing sphere volume with increasing dimensions!!!!!!



Objetos geométricos: la hiperesfera(II)

$$V(S^{15}) = \frac{\pi^8}{40320} R^{16} \approx 0,235331 R^{16} \quad (24)$$

$$V(S^{16}) = \frac{512\pi^8}{34459425} R^{17} \approx 0,140981 R^{17} \quad (25)$$

$$V(S^{23}) = \frac{\pi^{12}}{479001600} R^{24} \approx 0,00192957 R^{24} \quad (26)$$

$$V(S^{24}) = \frac{8192\pi^{12}}{7905853580625} R^{25} \approx 0,000957722 R^{25} \quad (27)$$

$$V(S^{25}) = \frac{\pi^{13}}{6227020800} R^{26} \approx 0,000466303 R^{26} \quad (28)$$

The amazing vanishing sphere volume with increasing dimensions!!!!!!



Objetos geométricos: la hiperesfera(II)

$$V(S^{26}) = \frac{16384\pi^{13}}{213458046676875} R^{27} \approx 0,000222872 R^{27} \quad (24)$$

$$V_{91} \left\{ \begin{array}{l} \frac{\pi^{46} R^{92}}{550262215981208894985030542880025489296165175296000000000000} \\ \approx 1,34377 \cdot 10^{-35} R^{92} \end{array} \right. \quad (25)$$

Two more... The 4096-dimensional sphere

$$V(S^{4095}) \approx 8,70008138919055 \times 10^{-4877} R^{4096} \quad (26)$$

with a fantastic fraction that can not be written in the margin or space of this page easily. Surprisingly, the infinite-dimensional sphere volume is zero:
The amazing vanishing sphere volume with increasing dimensions!!!!!!!!!

Objetos geométricos: la hiperesfera(II)

$$V(S^\infty) = 0 \quad (24)$$

The amazing vanishing sphere volume with increasing dimensions!!!!!!!!!

Bohr-logy in XD

2d Bohr energy levels and radius

In any 2d Universe, the Bohr-Rydberg energy and radius are:

$$E = K e^2 \left(\frac{1}{2} + \ln(n) \right), \quad r_n = n a_0 (2d) = \frac{n \hbar}{\sqrt{m K e}}, \quad n \in \mathbb{Z} \quad (25)$$

For gravitational case, take $K e^2 \rightarrow GMm$.

Dd Bohr energy levels and radius

In any Dd Universe, the Bohr-Rydberg energy and radius are:

$$E = \frac{D-4}{2(D-2)} \left(\frac{m}{\hbar^2} \right)^{\frac{D-2}{4-D}} n^{-\frac{2D-4}{4-D}} e^{\frac{4}{4-D}}, \quad r_n(D) = \left(\frac{m}{\hbar^2} \right)^{\frac{1}{D-4}} e^{\frac{2(2-D)}{(4-D)(D-2)}} n^{\frac{2}{4-D}} \quad (26)$$

For gravitational case, take $K e^2 \rightarrow GMm$.

Gravity in XD

Newton in higher dimensions

In any Dd ($D = d + 1$) Universe (spacetime), the gravitational force, the gravitational field, the potential energy and the potential read

$$F_N = G_D \frac{Mm}{r^{D-2}} = G_{d+1} \frac{Mm}{r^{d-1}} \quad g = G_D \frac{M}{r^{D-2}} = G_{d+1} \frac{M}{r^{d-1}} \quad (27)$$

$$U_g = G_D \frac{Mm}{r^{D-3}} = G_{d+1} \frac{Mm}{r^{d-2}}$$

$$V_g = G_D \frac{2\Gamma((D-1)/2)M}{\pi^{(D-3)/2} r^{D-3}} = G_{d+1} \frac{2\Gamma(d/2)M}{\pi^{(d-2)/2} (d-2) r^{d-2}} \quad (28)$$

Dilution of gravity: $G_N(4d) = G_D / V_{D-1}$. $g_{YM}^2(4d) = g_{YM,d}^2 R^{-d}$,

$M_P = \sqrt{\hbar c / G} \sim 10^{-5} g$, $M_W = \frac{\hbar}{c} \sqrt{\Lambda/3} \sim 10^{-65} g$. $G \hbar \Lambda / c^3 \sim 10^{-121}$.

$M_U = \frac{c^2}{G} \sqrt{3/\Lambda} \sim 10^{56} g$, $M'_W = \sqrt[3]{\frac{\hbar^2 \sqrt{\Lambda/3}}{G}} \sim 10^{-25} g$. $M_U / M_W \sim 10^{121}$

Uniform sphere total energy in XD

Gravitational/electric energy for uniform density sphere

$$U_g = -G_{d+1} \frac{d(d-2)M^2}{d+2} \frac{1}{R^{d-2}} = -G_D \frac{(D-1)(D-3)M^2}{D+1} \frac{1}{R^{D-3}} \quad (29)$$

with $D = d + 1$ and M the mass. If $M = \rho V$, then

$$U_g = -G_{d+1} \frac{d(d-2)\pi^d \rho^2}{(d+2)\Gamma^2(\frac{d}{2}+1)} R^{d+2} = -G_D \frac{(D-1)(D-3)\pi^d \rho^2}{(D+1)\Gamma^2(\frac{D+1}{2}+1)} R^{D+1} \quad (30)$$

Trickery for the electric case: substitute $G_n \rightarrow K_C$, $M \rightarrow Q$, with $Q = \rho V$.

Entropic gravity in XD

Hypothesis for $D = d + 1$ hyperdimensional Newton gravity:

- $A(\Sigma) = \frac{2\pi^{d/2} R^{d-1}}{\Gamma(d/2)}$.
- $N = A(\Sigma)/L_p^{d-1}$, $E = mc^2 = Nk_B T/2$, $\Delta S = 2\pi k_B \frac{mc\Delta x}{\hbar}$.

Then:

$$F = -T \frac{\Delta S}{\Delta x} = -G_d \frac{Mm}{R^{d-1}}$$

where

Hyperdimensional gravitational Newton constant

$$G_d = \frac{2\pi^{1-d/2} \Gamma\left(\frac{d}{2}\right) c^3 L_p^{d-1}}{\hbar} = 2\pi^{1-d/2} \Gamma\left(\frac{d}{2}\right) \frac{c^3 L_p^{d-1}}{\hbar}$$

$$\phi_g = -\Omega_d G_d M; \quad \phi_e = \Omega_d K_d Q = Q/\varepsilon_0(d) \quad \Omega_d = 2\pi^{d/2}/\Gamma(d/2)$$

Zeta function and gravitational constant

Take the functional equation:

$$\zeta(s) = 2^s \pi^{s-1} \sin\left(\frac{\pi s}{2}\right) \Gamma(1-s) \zeta(1-s) \text{ for } 1-s = d/2. \text{ Then, since}$$

$$G_d = 2\pi^{1-d/2} \Gamma\left(\frac{d}{2}\right) \frac{c^3 L_p^{d-1}}{\hbar}$$

you can derive that

Gravitational constant and zeta function

$$G_d = \frac{\pi 2^{d/2} \zeta\left(1 - \frac{d}{2}\right)}{\zeta\left(\frac{d}{2}\right) \cos\left(\frac{\pi d}{4}\right)} \left(\frac{c^3 L_p^{d-1}}{\hbar} \right) \quad (31)$$

Classical atom instability

Hypothesis: non quantum atoms are unstable. To prove this:

$$P = \frac{dE}{dt} = \frac{2e^2 a^2}{3c^3} \text{ (Larmor formula)}$$

$$\frac{Ke^2}{R^2} = \frac{mv^2}{R} \rightarrow v^2 = \frac{Ke^2}{mR}, E = \frac{mv^2}{2} - \frac{Ke^2}{R} = -\frac{Ke^2}{R}$$

$$dt = -\frac{1}{dE} \frac{dE}{dR} dR = -\frac{3}{16} \frac{m^2 c^3 R^2 dR}{(E_0 R_0)^2} \rightarrow \int_0^{t_c} dt = -\frac{3m^2 c^3}{(E_0 R_0)^2} \int_{R_0}^0 R^2 dR$$

We finally get:

Decay time of classical em-atoms

$$t_c = \frac{m^2 c^3 R_0}{16 E_0^2} = \frac{m^2 c^3 R_0^3}{4 K_C e^4} = \frac{4\pi^2 \varepsilon_0^2 m^2 c^3 R_0^3}{e^4} \simeq 1,6 \cdot 10^{-11} s \sim 20 ps$$

Gravitational music equations

Binary system with $M = M_1 + M_2$, $f_{GW} = 2f_{orb}$, $M_c = (M_1 M_2)^{3/5}/M^{1/5}$ yields (GR):

$$L_{GW} = \frac{2^5}{5} \left(\frac{G^{7/3}}{c^5} \right) [M_c \pi f_{GW}]^{10/3}$$

$$\dot{f}_{GW} = \left(\frac{96}{5} \right) \left(\frac{G^{5/3}}{c^5} \right) \left(\pi^{8/3} \right) (f_{GW})^{11/3}$$

$$t_c = \frac{2}{2^8} \left(\frac{GM_c}{c^3} \right)^{-5/3} [\pi f_{GW}]^{-8/3}$$

Neutrino oscillations: the equations

Supposing transitions between different neutrino species, via $|\nu_\alpha\rangle$ to $|\nu_\beta\rangle$

Neutrino oscillations

$$\mathcal{A} = P(\alpha \rightarrow \beta) = |\langle \nu_\beta | \nu_\alpha \rangle|^2 = \left| \sum_j U_{\alpha j}^* U_{\beta j} e^{-i \frac{m_j^2 c^3 L}{2\hbar E}} \right|^2 \quad (32)$$

$$\mathcal{A} = \delta_{\alpha\beta} - 4 \sum_{i>j} \operatorname{Re} (U_{\alpha i}^* U_{\beta i} U_{\beta j} U_{\beta j}^*) \sin^2 \left(\frac{\Delta m_{ij}^2 c^3 L}{4\hbar E} \right) +$$

$$+ 2 \sum_{i>j} \operatorname{Im} (U_{\alpha i}^* U_{\beta i} U_{\beta j} U_{\beta j}^*) \sin \left(\frac{\Delta m_{ij}^2 c^3 L}{2\hbar E} \right) \quad (33) \text{ with}$$

Neutrino transition matrix

$$|\nu_\beta\rangle = U_{\beta}^\alpha |\nu_\alpha\rangle$$

Multitemporal physics(I)

Usual 1T newtonian physics: $F = ma = m\frac{dv}{dt} = m\frac{d^2r}{dt^2} = -\nabla U(r)$,
assuming conservative forces only. Let $W = F_i dx^i$ the work form, in a ND
manifold $V \subset \mathbb{R}^N$, with submanifold nd $M \subset \mathbb{R}^n \subset \mathbb{R}^N$. $y^I = y^I(x)$,

$\omega = f_I dy^I$ implies $dy^I = \frac{\partial y^I}{\partial x^i} dx^i$, and also

$$W = F_i(x) dx^i \rightarrow F_I = f_I(y(x)) \frac{\partial y^I}{\partial x^i}$$

Single time manifold approach

$$f_I = m\delta_{IJ} \frac{d\dot{y}^J}{dt} = m\delta_{IJ} \frac{d^2y^I}{dt^2}$$

$$F_i = m\delta_{IJ} \frac{d\dot{y}^I}{dt} \frac{\partial y^J}{\partial x^i} = m\delta_{IJ} \frac{d^2y^I}{dt^2} \frac{\partial y^J}{\partial x^i}$$

Multitemporal physics(II)

Going multitemporal with timelike coordinates $(t) = t^\alpha$, $\alpha = 1, \dots, m$

Multitime tensorial Newton 2nd law

$$f_I = m_{IJ} \delta^{\alpha\beta} \frac{\partial^2 y^J}{\partial t^\alpha \partial t^\beta}$$

$$f_i = m_{IJ} \delta^{\alpha\beta} \frac{\partial^2 y^I}{\partial t^\alpha \partial t^\beta} \frac{\partial y^J}{\partial x^i}$$

with anti-trace $F_i = F_{i\alpha}^\alpha$ given by the tensor 1-form

$$F_{i\alpha}^\sigma = m_{IJ} \delta^{\sigma\beta} \frac{\partial^2 y^I}{\partial t^\alpha \partial t^\beta} \frac{\partial y^J}{\partial x^i}$$

(Multitime) Kinetic energy

$$T = E_k = \frac{1}{2} m \delta_{IJ} \dot{y}^I \dot{y}^J \quad T = \frac{1}{2} \delta_{IJ} \delta^{\alpha\beta} \frac{\partial y^I}{\partial t^\alpha} \frac{\partial y^J}{\partial t^\beta}$$

Multitemporal physics(III)

Single time Euler-Lagrange 1st order EOM

$$\delta S = 0 \rightarrow E(L) = \frac{\partial L}{\partial x^i} - \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}^i} \right) = 0$$

(Multitime) Euler-Lagrange EOM

$$\delta S = 0 \rightarrow E(L) = \frac{\partial L}{\partial x^i} - D_\alpha \left(\frac{\partial L}{\partial D_\alpha x^i} \right) = 0$$

(Multitime) Euler-Lagrange EOM: nth order

$$E(L) = \sum_{j=0}^n (-1)^j \left(\frac{\partial L}{\partial \partial_t^j x^i} \right) = 0 \quad E(L) = \sum_{J=0}^n (-1)^J \left(\frac{\partial^J L}{\partial D_\alpha^J x^i} \right) = 0$$

Multitemporal physics(IV)

Single time Hamilton EOM

Define 1T hamiltonian as $H = \dot{x}^i \frac{\partial L}{\partial \dot{x}^i} - L$, and $p_i = \partial L / \partial \dot{x}^i$, then

$$\dot{x}^i = \frac{dx^i}{dt} = \frac{\partial H}{\partial p_i} \quad \dot{p}_i = \frac{dp_i}{dt} = -\frac{\partial H}{\partial x^i}$$

Multi-time Hamilton EOM

Define nT hamiltonian as $H = D_\alpha x^i \frac{\partial L}{\partial D_\alpha x^i} - L$, and $p_i^\alpha = \partial L / \partial D_\alpha x^i$, then

$$\frac{\partial x^i}{\partial t^\alpha} = \frac{\partial H}{\partial p_i^\alpha} \quad \frac{\partial p_i^\beta}{\partial t^\alpha} = -\delta^\beta_\alpha \frac{\partial H}{\partial x^i}$$

Detecting exoplanets(I)

Astrometry

$$\theta = \left(\frac{M_p}{M_\star} \right) \left(\frac{a}{r} \right) \approx \frac{10^{-3}}{r(\text{pc})} \left(\frac{P(\text{yr})}{M_\star(\odot)} \right)^{2/3} M_p(J)$$

Here

$$V_r(m/s) \approx \frac{30}{(P(\text{yr}))^{1/3}} \frac{M_p(J)}{M_\star(\odot)^{2/3}} \sin(i)$$

Detecting exoplanets(II)

Microlensing

$$R_E^2 = \frac{4GMD}{c^2}, \quad D = \frac{D_{ds}D_d}{D_s}, \quad t_0 = \frac{R_E}{v_e}$$

$$t_0 = \frac{2D_L\theta_E}{v_L} = \frac{2\theta_L}{v_L} \sqrt{\frac{4GM(1 - D_s/D_s)}{c^2 D_d}}$$

The impact parameter u reads

$$A = \frac{u^2 + 2}{u(u^2 + 4)^{1/2}}$$

Detecting exoplanets(III)

Direct detection

$$B \geq \frac{\lambda D}{r} \approx \left(\frac{\lambda}{10\mu mu} \right) \left(\frac{D}{10pc} \right) \left(\frac{r}{1AU} \right)^{-1} m$$

Detecting exoplanets(IV)

Radial velocity

$$K_{\star} = \left(\frac{2\pi G_N}{P} \right)^{1/3} \frac{M_p (M_{\star} + M_p)^{1/3} \sin(i)}{M_{\star}} \frac{1}{\sqrt{1 - e^2}}$$

Also, it is usually written with $M_{\star} + M_p \simeq M_{\star}$ as follows

$$M_p \sin(i) = \left(\frac{P}{2\pi G} \right)^{1/3} K_{\star} M_{\star}^{2/3} \sqrt{1 - e^2}$$

Bohr-like quantization of magnetic monopoles

Hypothesis:

- Magnetic and electric field of a point monopole charge with $Q_m = g$ and dual charge $e_g = eg/c = egv/c^2$.

$$F_e + F_m = 2F_{m,e} = F_c \leftrightarrow \frac{2K_C e_g}{R^2} = \frac{mv^2}{R} \rightarrow \frac{c^{-2}eg}{4\pi\varepsilon_0} = \frac{mvR}{2} = \frac{n\hbar}{2}$$

Then, $eg = \frac{n\hbar c^2}{2K_C}$ (Q.E.D.). Equivalently: $\frac{g}{e} = \frac{nc}{2\alpha_e} \leftrightarrow \alpha_e = \frac{nce}{2g}$

- Dirac-Zwanziger-Schwinger dyonic quantization $Z = (e, g)$:

$$e_1g_2 - e_2g_1 = 2\pi n\hbar c$$

From this, it follows that $Q = ne$, or $Q = \left(n + \frac{1}{2}\right)e$ and

$$M = \sqrt{\frac{K_C}{G_N}}e = \frac{\hbar c^2}{g} \sqrt{\frac{1}{K_C G_N}}.$$

The existence of magnetic monopoles implies the quantization of Q_e .



Doctor Strange in the Multiverse of Madness!

